

# A Nanopore–Nanofiber Mesh Biosensor To Control DNA Translocation

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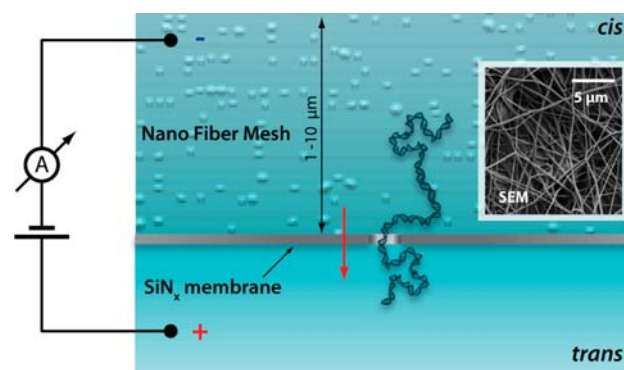
**S** Supporting Information

**ABSTRACT:** Solid-state nanopores show promise as single-molecule sensors for biomedical applications, but to increase their resolution and efficiency, analyte molecules must remain longer in the nanopore sensing volume. Here we demonstrate a novel, facile, and customizable nanopore sensor modification that reduces the double-stranded DNA translocation velocity by 2 orders of magnitude or more via interactions outside the nanopore. This is achieved by electrospinning a copolymer nanofiber mesh (NFM) directly onto a solid-state nanopore (NP) chip. The effect of NFMs on dsDNA translocation through an NP is highlighted using a set of NFMs of varying mesh composition that reduce the translocation speed relative to a bare pore from 1- to >100-fold. A representative NFM from this set is effective on DNA as long as 20 kbp, improves the nanopore resolution, and allows discrimination among different DNA lengths.

The emergence of novel nanomaterials and nanofabrication tools is accelerating the development of single-molecule biosensors to directly detect genomic information, such as single-nucleotide polymorphisms, structural variations, and epigenetic markers. One of the simplest and most versatile biosensors in this class is the solid-state nanopore (NP).<sup>1</sup> Similar to protein-based NPs,<sup>2</sup> solid-state NPs employ electrophoretic forces to thread and slide electrically charged biopolymers through a nanoscale hole fabricated in an ultrathin insulating membrane. A large voltage potential applied across an NP immersed in an electrolyte solution induces an intense ion current, producing a strongly diverging electric field in the vicinity of the pore. This field gradient efficiently focuses electrically charged biopolymers, such as DNA molecules, allowing the detection of minute sample concentrations.<sup>3</sup> The potential ability of an NP to scan and to resolve small, local features along DNA necessitates that the nominal thickness of the membrane be limited to only a few nanometers or even less.<sup>4</sup> However, such thin membranes significantly limit the surface of the pore available to interact with DNA and slow its movement through the NP. Typical translocation (or sliding) speeds of double-stranded DNA (dsDNA) in solid-state NP range from tens to hundreds of nanoseconds per base pair (bp)<sup>1a</sup> and cannot be resolved with sufficient accuracy because of the inherent electrical noise in the detection system at these bandwidths.<sup>5</sup> Thus, methods capable of slowing (or better yet controlling) translocation speeds are of interest for sensing

applications such as DNA and RNA sequencing, gene expression, and genotyping.<sup>6</sup>

To regulate the translocation speed of DNA molecules, we have developed a highly porous and tunable synthetic coating that is applied to either the entry or the exit (cis or trans, respectively) faces of a thin silicon nitride membrane containing the pore, as depicted schematically in Figure 1.



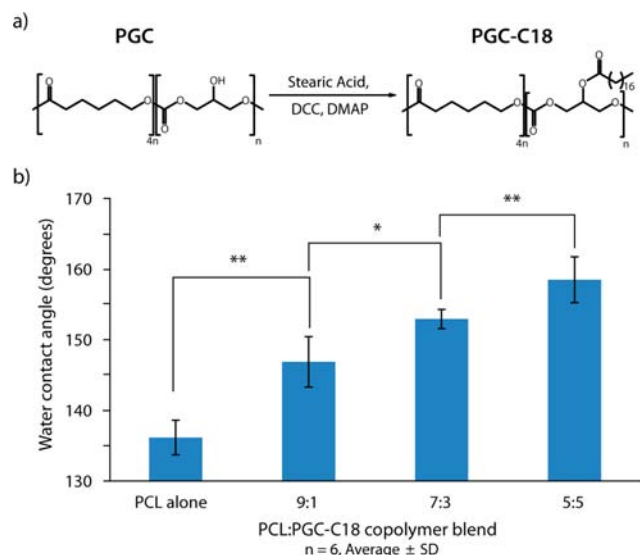
**Figure 1.** Schematic cross section of a solid-state nanopore with a nanofiber mesh spun on the cis side. The DNA interacts with the mesh as it is electrophoretically threaded through the pore (not to scale). Inset: SEM image of a nanofiber mesh on a nanopore chip.

The coating is formed by a polymeric nanofiber mesh (NFM) with tunable chemical and physical properties that is electrospun directly onto a surface.<sup>7</sup> Previous studies have demonstrated that NFMs can act as an adsorbent to separate biopolymers.<sup>8</sup> We propose that a low-density, high-surface-area NFM proximal to an NP will significantly slow DNA passage by interacting with the biopolymer prior to and during translocation through the NP. Our approach differs from previous strategies for NP modification, which influence translocation by altering the pore surface itself,<sup>9</sup> directly tethering and manipulating the DNA,<sup>10</sup> or changing properties of the surrounding medium such as viscosity, pressure, or ionic strength.<sup>11</sup> These approaches typically do not decouple improvements in translocation dynamics from other characteristics of the device, such as conductivity, blockage level, or wall charge. This can lead to undesirable consequences, including reduced threading efficiency and ion current stability, smaller signal-to-noise ratios, and reduced NP hydration efficiency.

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NFMs were formulated from copolymer blends of poly( $\epsilon$ -caprolactone) (PCL) (70–90 kg/mol, Sigma) and poly(glycerol monostearate-*co*- $\epsilon$ -caprolactone) (PGC-C18) (22 kg/mol) (Figure 2a). PGC-C18 was synthesized according to

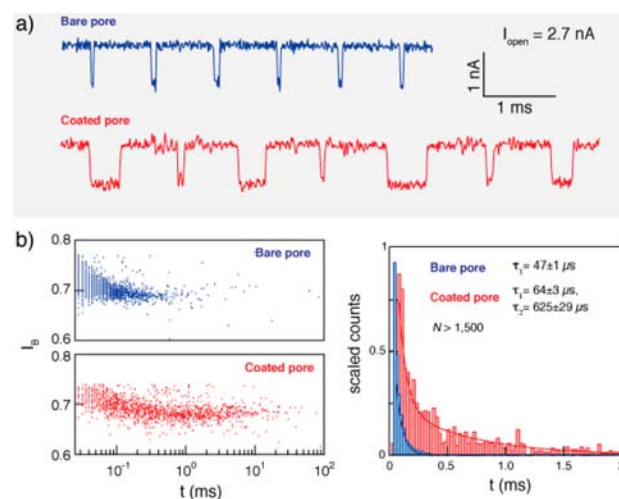


**Figure 2.** (a) Synthesis of PGC-C18. (b) Effect of the PCL:PGC-C18 copolymer ratio on the nanofiber mesh hydrophobicity, as measured via the contact angle ( $n = 6$ ; \*,  $p < 0.05$ ; \*\*,  $p < 0.01$ ).

our previously published protocol [see the Supporting Information (SI)].<sup>7a,12</sup> Doping PCL with increasing quantities of the hydrophobic PGC-C18 increases the resulting mesh hydrophobicity, as characterized by the water droplet contact angle (Figure 2b). These measurements of hydrophobicity indicate changes in chemical composition. The Figure 1 inset shows a typical scanning electron microscopy (SEM) image of a 7:3 PCL:PGC-C18 hydrophobic NFM electrospun onto an NP chip, with fiber diameters ranging from 300 to 450 nm. In all cases, the electrospinning deposition time, voltage, and needle position were adjusted to produce uniform NFM thicknesses with fibers of similar morphology across all polymer blends (see Table S1 and Figures S5–S7 in the SI). The NFM fabrication step is facile, parallel, and fast. For example, coating 50 chips required about 1 min and a few milligrams of polymer. The resulting NFM has a low volume fraction and high surface area and is mechanically stable. In this study, we electrospun NFMs onto NP chips using the following PCL:PGC-C18 copolymer blend ratios: 10:0 (PCL only), 9:1, 8:2, 7:3, 6:4, and 5:5.

The nanopore–nanofiber mesh (NP–NFM) sensor consists of a small NP drilled with a tightly focused transmission electron microscope electron beam in an LPCVD-deposited, low-stress  $\text{SiN}_x$  membrane (25 nm thick; Figure S1).<sup>13</sup> NP chips were sealed in a custom-built flow cell permitting a low-noise recording of the ion current flowing through the pore.<sup>13</sup> DNA added to the cis side of an unmodified NP (no NFM) under an applied electric potential induces blockades in the ionic current corresponding to translocations of DNA from the grounded cis side to the positively biased trans side of the membrane. Typical translocation events for 1000 bp DNA in an unmodified bare  $\text{SiN}_x$  NP are shown in Figure 3a (blue).

A 7:3 PCL:PGC-C18 polymer solution was electrospun onto the very same NP. The NP–NFM device was readily hydrated and permitted both buffer and sample exchange (see Figure S8). Translocations of 1000 bp DNA through this NP–NFM



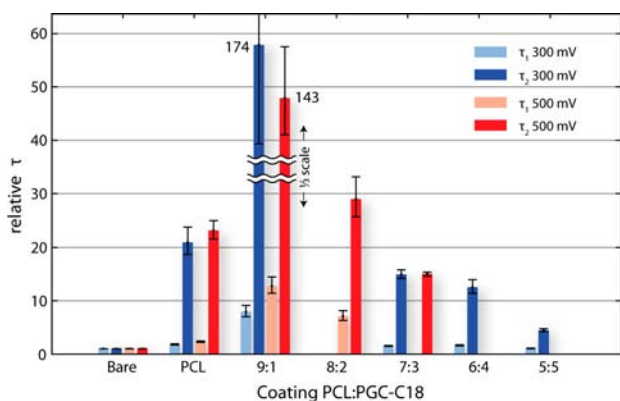
**Figure 3.** Comparison of translocation times through a bare nanopore (NP) and an NP–NFM. (a) Representative current traces for 1000 bp dsDNA passing through the same NP as a bare NP (upper) and as a 7:3 PCL:PGC-C18 NP–NFM (lower), showing similar blockage levels but some significantly longer translocation times for the NP–NFM ( $V = 300$  mV,  $I_{\text{open}} = 2.7$  nA). (b) Histogram of event duration for the same uncoated NP (blue) and 7:3 PCL:PGC-C18 NP–NFM (red), showing exponential tail fits (error bars:  $\tau \pm 95\%$  confidence interval).

(Figure 3) revealed two important characteristics of the modified NP: First, and most noticeably, we observed a broader spread in the DNA translocation time. Specifically, the dwell time of a large fraction of the events falls between 0.5 and 10 ms (Figure 3b), a range that exceeds the typical translocation time for the same uncoated pore by roughly an order of magnitude. Second, the presence of the NFM does not substantially affect the open pore current (the ion current prior to DNA entry into the pore), the blocked current level, or the noise in the NP (Figure S9). A closer evaluation of the translocation events (see the sample events in Figure 3a and the dwell time histograms in Figure 3b) suggests that instead of a uniform shift of the entire dwell time histogram toward longer timescales, the NFM induces a bimodal distribution containing populations of “normal” and “long” events. Indeed, a monoexponential tail fit failed to represent the dwell time histogram of the NFM-coated pore as accurately as a double-exponential fit. The shorter timescale,  $\tau_1$ , is close to the typical timescale for the uncoated pore, while the longer timescale,  $\tau_2$ , is nearly 10 times longer. A detailed analysis of the fits and errors obtained for all data sets are presented in Table S2.

To further characterize the nature of the slowed DNA translocations induced by the presence of a tunable NFM, we fabricated and tested a set of NP–NFMs representing six different PCL:PGC-C18 copolymer blends spun onto 4–4.5 nm pores. Similar to our observations for the 7:3 copolymer blend, the relative blockage level ( $I_{\text{block}} = I_{\text{block}}/I_{\text{open}}$ ) and the conductance measured were nearly the same across all of these compositions (Figures S10 and S11). These results suggest that the ion mobility near the NP is similar to that of a bare NP, consistent with a highly porous NFM structure.

In contrast to the ion current levels, the dsDNA translocation dynamics were highly dependent on the NFM composition. We measured the characteristic translocation times of 1000 bp dsDNA using different NFM copolymer blend coatings, once again tail-fitting the resulting dwell time distributions to double-

exponential functions (see Figures S14 and S15 for all of the fits). Other metrics for analyzing the translocation time were explored [e.g., histogram peaks ( $t_p$ ) for logarithmic binning of the data], but these approaches generally failed to capture both the short and long translocation event populations (see Table S2 and Figure S13). We defined the relative  $\tau$  ( $\tau_{\text{relative}}$ ) as the ratio of the timescale for the “long” event population ( $\tau_2$ ) at each coating normalized by the characteristic timescale of translocation for the bare pore. We repeated these measurements at two applied voltages (300 and 500 mV). Our results are summarized in Figure 4 (error bars show corresponding



**Figure 4.** Slowing factor  $\tau_{\text{relative}}$  ( $=\tau_{\text{coated}}/\tau_{\text{bare}}$ ) for various coatings: Bare pore ( $\tau_{\text{relative}} = 1$ ), PCL only, and 9:1, 8:2, 7:3, 6:4, and 5:5 PCL:PGC-C18 blends. All data are for 1000 bp dsDNA in 4–4.5 nm nanopores at 300 mV (blue) or 500 mV (red) (error bars:  $\tau \pm 95\%$  fit confidence interval of exponential tail fits; the 9:1 PCL:PGC-C18 blend data at 300 and 500 mV are shown at 1/3 scale for clarity).  $\tau_{\text{relative}}$  values calculated using  $\tau_{\text{coated}} = \tau_1$  or  $\tau_2$  (light or dark colors, respectively) are shown.

95% confidence intervals for the fits). For reference, we also show the relative short  $\tau$  ( $\tau_{\text{rel\_short}}$ ) values using the normal translocation population ( $\tau_1$ ) where available;  $\tau_{\text{rel\_short}}$  generally showed values around unity.

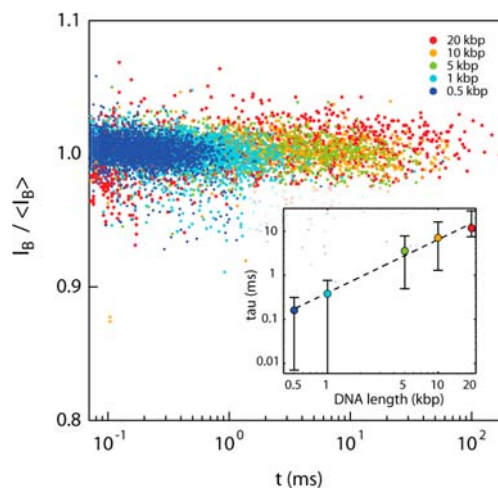
When the NFMs were ranked in order of increasing hydrophobicity according to contact angle measurements, we observed non-monotonic changes in  $\tau_{\text{relative}}$ : The most- and least-hydrophobic NFMs, respectively, had relatively little effect on the translocation speed. PCL alone slowed translocations by more than 20-fold at both 300 and 500 mV. The super-hydrophobic 6:4 and 5:5 PCL:PGC-C18 meshes slowed DNA by only 12- and 4-fold, respectively, at the lower driving force of 300 mV. For intermediate copolymer blends, the data collected at both 300 and 500 mV clearly showed a more pronounced slowing effect than the most- and least-hydrophobic meshes. In particular, the 9:1 PCL:PGC-C18 NFM slowed translocations by >140-fold at 500 mV and >170-fold at 300 mV. At 300 mV, nearly 20% of events for this mesh were longer than 10 ms. For comparison, <0.2% of events in the bare pore at 300 mV were longer than 10 ms.

The variation of  $\tau_{\text{relative}}$  with mesh composition suggests that the translocating DNA interacts with the NFM as it approaches and threads through the NP and that the strength of these interactions changes with the chemical composition of the mesh. Moreover, the fact that the values obtained for  $\tau_1$  are close to the bare pore translocation times suggests that only a fraction of the DNA molecules interact with the NFM. This observation is consistent with the presence of very sparse

NFMs. The interactions of strong polyelectrolytes, such as DNA, with dielectric surfaces are governed by a complex interplay between electrostatic and hydrophobic forces, which depend not only upon chemical composition but also on the material’s structure and texture and other steric considerations.<sup>14</sup> The most hydrophobic NFMs produce relatively small retardation effects, while the NFMs characterized by intermediate hydrophobicity levels create maximum drag on the DNA. To account for this complex behavior, a detailed model describing the DNA–NFM interactions must be developed. From a practical standpoint, however, this interplay provides flexibility in tuning the material properties of the NFMs.

Finally, we collected translocation events using the 7:3 PCL:PGC-C18 copolymer blend NFM at 500 mV for five different dsDNA lengths ranging from 0.5 to 20 kbp to determine whether longer biopolymers interact more strongly with the NFM than shorter biopolymers. One might expect that the number of contact points between the mesh and DNA would increase with biopolymer length, affording a more stable overall interaction. To maintain consistency across the samples, all of the measurements were performed sequentially in a single 6 nm diameter pore with the same NFM coating, and some data sets were collected twice at different time points to ensure reproducibility. The characteristic ion current level and dwell time of each event was extracted and plotted on an “event diagram” (Figure 5). Events that displayed a folded DNA translocation pattern<sup>15</sup> were excluded in the analysis to simplify interpretation of the results.

Figure 5 shows a clear pattern of longer translocation times with larger DNA molecules. While we did not make an attempt



**Figure 5.** Event diagram for five lengths of dsDNA translocating through a 6 nm nanopore coated with 7:3 PCL:PGC-C18.  $I_b$  is normalized for clarity. Inset: plot of characteristic translocation time  $\tau$  vs DNA length (error bar:  $\tau \pm 95\%$  fit confidence interval). The dotted line is a guide to the eye.

to discriminate collision events (fast events that involve unsuccessful threading of the DNA into the pore) from true translocations, the overall trend of the translocation time is clear and consistent for all lengths. As before, we numerically characterized the translocation dwell time distributions using exponential tail fits (see Figures S16 and S17). These results, shown in the inset of Figure 5, indicate mean translocation speeds of roughly 0.4–0.7  $\mu\text{s}/\text{base}$ , which are 20–35-fold



slower than for an uncoated pore under the same conditions (see Figure S18). A monotonic growth in the characteristic translocation time as a function of length was observed for DNA in the presence of the NFM coating. We also observed that in the range from 1 to 10 kbp, the slowing factor relative to a bare pore increased slightly (from 20 to 35; see the Figure S18 inset). While this is consistent with our original hypothesis, the trend of increased slowing for longer DNA was far less pronounced than expected and barely significant given the associated fit error. Although this observation partly contradicts our a priori expectation that the longest DNA would be slowed much more than shorter DNA, there are still a number of possible explanations for this behavior. First, some of the events in the 20 kbp sample and even the 10 kbp sample exceeded the acquisition capability of our experimental system (~250 ms); thus, the overall tail fit may reflect shorter timescales than expected. Second, a fully stretched 20 kbp DNA may extend beyond the width (even locally) of the NFM fibers used in this experiment. It is thus reasonable to predict that the retardation factor may stay constant or even become smaller for very long DNAs. Nevertheless, a clear relationship between the characteristic translocation time and DNA length exists.

In summary, the effect of NFM coatings on dsDNA translocation dynamics in solid-state NPs is reported. The NFMs increase DNA translocation time by up to 2 orders of magnitude or more without altering the ion current levels. This effect is sustained for DNA up to 20 kbp in length, enabling greater temporal resolution for the longest strands of DNA. The NFM composition, as characterized by hydrophobicity, affects the translocation time. This observation is consistent with our view that at an intermediate hydrophobicity the DNA interacts strongly with the NFM, the disruption of which is facilitated by the electrophoretic forces applied on the DNA. The process of electrospinning an NFM coating onto an NP is facile, high-throughput, parallel, and compatible for use with a number of chemically diverse polymers. Thus, this method and the resulting device compositions can be readily adjusted for many DNA-sensing applications benefiting from control over biopolymer translocation rates. Future work will focus on these and other NFMs to create an NP-based class of biosensors with a broad range of customizable translocation properties.

## ■ ASSOCIATED CONTENT

### ● Supporting Information

Experimental procedures, synthesis, characterization, statistics, and supporting figures. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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### Notes

The authors declare no competing financial interest.

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